



# Highly compatible and reliable ZrN interfacial layer between TiN top electrode and antiferroelectric ZrO<sub>2</sub> thin film to boost the electrocaloric behavior

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## ABSTRACT

A highly compatible and reliable ZrN interfacial layer with different cycles of atomic layer deposition was inserted between the TiN top electrode and antiferroelectric (AFE) ZrO<sub>2</sub> thin film to achieve low oxygen vacancies and high crystallinity of the ZrO<sub>2</sub> thin film, boosting the AFE and electrocaloric (EC) characteristics of TiN/ZrO<sub>2</sub>/TiN metal–insulator–metal (MIM) capacitors. An outstanding adiabatic temperature change ( $\Delta T$ ) of  $-21$  K was realized for devices with a ZrN interfacial layer of nine cycles, which can be ascribed to the reduction in redox between the TiN and ZrO<sub>2</sub> for the suppression of interfacial dead layer during annealing as confirmed via high-angle annular dark field images and interfacial capacitance measurement. Hence, the reliability of ZrO<sub>2</sub> MIM capacitors with a nine-cycle ZrN interfacial layer was examined with a negligible change in AFE and EC behaviors after a cycling endurance test of  $10^6$  cycles, exhibiting great potential for future applications in solid-state cooling.

## 1. Introduction

With a rapid development of electronics industry to explore a metaverse, the high-density integrated circuits (ICs) with a fast clock rate are inevitable and often accompanied by a severe thermal design power (TDP) [1]. It is predicted that the TDP of high-performance graphic processing units (GPUs) will reach at least 700 watts before 2025 [2]. Hence, the demand for efficient cooling systems in ICs has increased urgently. However, the application of traditional cooling fan and heat sink in ICs encounters some bottlenecks such as the low efficiency, scaling limitation, and environmental pollutions [3]. Thus, some techniques using the electrocaloric (EC) effect, pyroelectric effect, or magnetocaloric effect have been proposed for the next-generation cooling in ICs. Among them, EC effect has been considered as the most promising method because of its mature manufacturing in complementary metal-oxide-semiconductor (CMOS), good scalability, and simple

operation [4]. Electrocaloricity is a phenomenon which can be defined as the adiabatic temperature change ( $\Delta T$ ) induced by applying or removing an external electric field in dielectric materials with the polarized behaviors [5]. When an external electric field is applied on the dielectric material, the originally chaotic dipole will be aligned with the direction of the field. In an adiabatic system, to compensate the entropy disturbance caused by the polarization, the material would generate additional entropy as offset, which would be manifested in the form of temperature to warm up or cool down the system [6]. The  $\Delta T$  can be quantified as a temperature change with the electric field and directly measured via the thermal microscopy, thermocouples, and infrared camera [7]. Besides, it can be indirectly extracted from the polarization versus temperature curves under different electric fields by using Maxwell's equations [8].

Additionally, the EC effect has been reported in various dielectric materials, including polymers, perovskite ceramics, and high dielectric

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constant (high- $k$ ) ceramic thin films [9–11]. Among these materials, a large  $\Delta T$  has been achieved in ferroelectric hafnium zirconium oxide (HZO) and antiferroelectric (AFE)  $ZrO_2$  thin films via an atomic layer deposition (ALD) system with a simple process, high scalability, and robust reliability [12,13]. Compared to ferroelectric materials, the AFE materials exhibit double polarized hysteresis loops with zero remnant polarization at zero bias, which can be ascribed to the field-induced phase transition between the non-polar tetragonal phase to the polar orthorhombic phase [14]. Therefore, AFE materials can demonstrate stable polarization switching and high fatigue resistance due to their relatively small depolarization field, suitable for the EC application in ICs [15]. Nevertheless, there are still some issues at the interface between the high- $k$  ceramic layer and metal electrode such as the redistribution of oxygen vacancy, electrode oxidation, and plasma damage, leading to the extra defects and unwanted interfacial layer for the degraded polarization and EC characteristics [16,17]. Previously, methods have been proposed to solve the interfacial problems at bottom electrode (BE) interface of ferroelectric HZO capacitors, which are originated from the oxidation of TiN during the ALD deposition. HfN and in-situ TiN have been inserted between HZO and TiN layers, giving rise to the inhibition in the formation of interfacial  $TiO_x$  dead layer [18]. Additionally, Chen et al. presented that the redistribution of oxygen vacancy at the interface can be effectively suppressed via the  $NH_3$  plasma surface treatment [19]. Thus far, these technologies have been applied at the BE interface of ferroelectric HZO capacitors to improve their memory characteristics. In this study, the enhancement of AFE and EC behaviors in TiN/ $ZrO_2$ /TiN metal–insulator–metal (MIM) capacitors via the interfacial zirconium nitride (ZrN) passivation at the top electrode (TE) interface has been proposed, which is relatively more compatible for the  $ZrO_2$  thin films and has not yet been discussed before. ZrN has excellent corrosion and oxidation resistance properties, which has been used as diffusion barriers [20,21], oxygen reduction catalysts [22], and biomedical coatings [23,24]. Here, the influence of a ZrN interfacial layer with different ALD cycles on the redox reaction between the TiN and  $ZrO_2$  during the annealing process as well as the AFE and EC characteristics and reliabilities of the AFE  $ZrO_2$  thin film was investigated. The enhanced behaviors of the TiN/ $ZrO_2$ /TiN MIM capacitors with a ZrN interfacial layer can be ascribed to the thin interfacial dead layer, resulted from the less redox reaction. Furthermore, the robust reliability under the cycling endurance tests was realized because of the suppression of the oxygen vacancy generation during the cycling operation.

## 2. Experimental

### 2.1. Sample preparation

The AFE  $ZrO_2$  capacitors with a ZrN interfacial layer were fabricated on 4-inch  $n^+$ -silicon wafers. After the removal of the native oxide using a diluted HF solution, a 100-nm-thick TiN film was deposited via a dc sputtering system (G3.5 Cluster, FSE Corp., New Taipei City, Taiwan) at 5 kW as the bottom electrode (BE). The gas flow rates of argon and nitrogen were kept at 30 and 30 sccm, respectively, to maintain a working pressure of 3 mTorr. Then, a 10-nm-thick  $ZrO_2$  layer was deposited via a plasma-enhanced ALD (PEALD) system (Cambridge NanoTech Fiji-202 DCS, Oxford Instruments, Oxford, UK) at 250 °C. Subsequently, a ZrN interfacial layer with different ALD cycles of 3, 6, 9, and 12 were deposited on the  $ZrO_2$  layer and denoted as ZrN\_3, ZrN\_6, ZrN\_9, and ZrN\_12, respectively. A reference device without a ZrN interfacial layer was also prepared and denoted as w/o. For the deposition of  $ZrO_2$  and ZrN layers, tetrakis(dimethylamino) zirconium (TDMAZ),  $O_2$  plasma, and  $NH_3$  plasma were used as the precursors of Zr, O, and N, respectively. After that, a 100-nm-thick TiN film was deposited on all devices as the top electrode (TE) and then annealed at 600 °C for 30 s in  $N_2$  ambient (AccuThermo AW 810 RTP, Allwin21 Corp., Morgan Hill, CA, USA) to form the AFE phase of the  $ZrO_2$  thin films. Finally, the TE was

patterned via photolithography and then etched using a  $NH_4OH:H_2O_2:H_2O$  (1:1:5) solution to obtain the isolated electrodes with a diameter of 100  $\mu m$ .

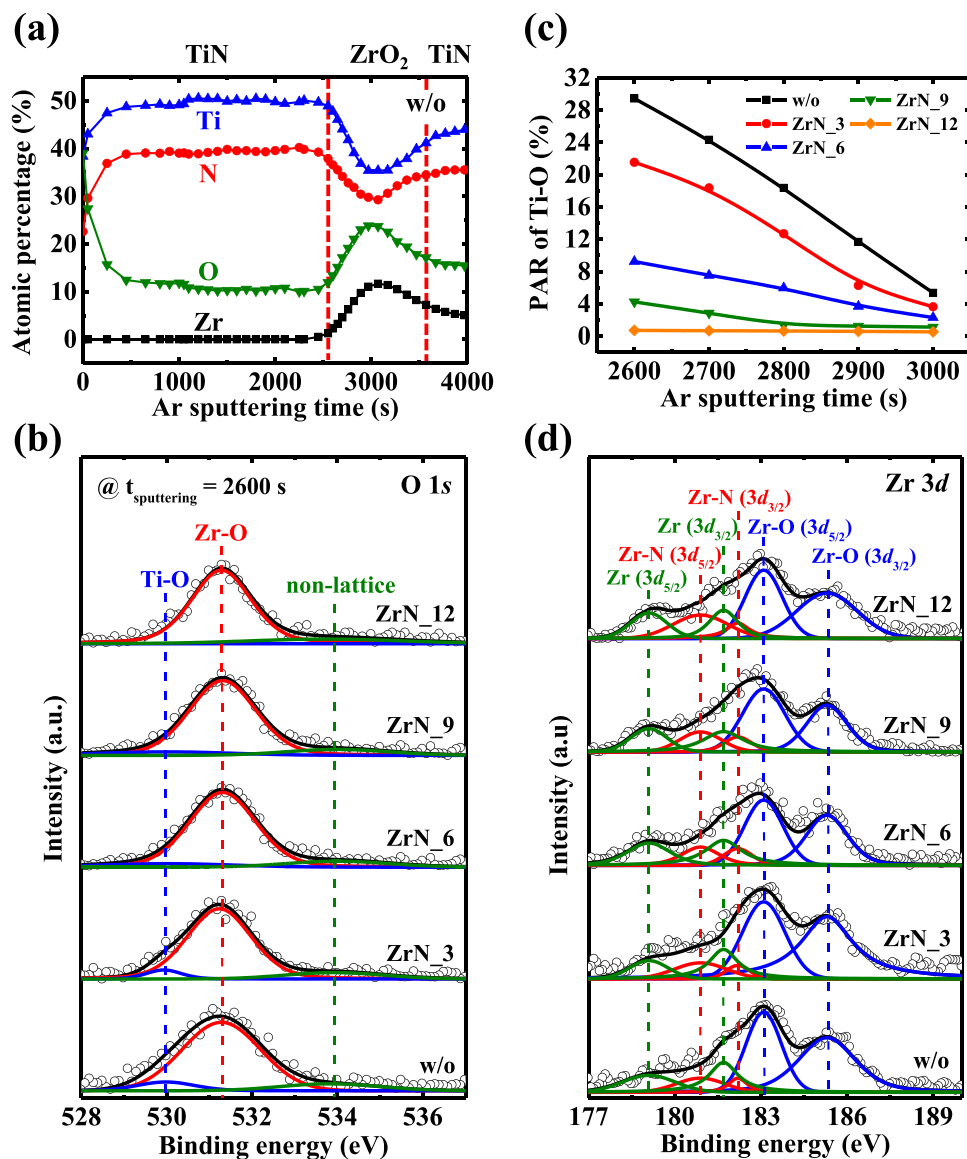
### 2.2. Characterization of materials and devices

The chemical composition at the interface between the TiN TE and  $ZrO_2$  thin film was analyzed using X-ray photoelectron spectroscopy (XPS) (Theta Probe ARXPS, Waltham, KS, USA) at  $10^{-9}$  Torr with Al  $K\alpha$  as the X-ray source (1486.7 eV). Furthermore, the crystallization behaviors of  $ZrO_2$  thin films were confirmed via grazing incidence X-ray diffraction (GIXRD) (X'Pert Pro MRD, Almelo, Netherlands) with Cu  $K\alpha$  radiation ( $\lambda = 0.154$  nm) at an incident angle of  $0.5^\circ$  and high-resolution transmission electron microscope (HRTEM) (JEM-2100PLUS, JEOL, Tokyo, Japan). To investigate the polarized phases of the TiN/ $ZrO_2$ /TiN stacked films with a ZrN interfacial layer, high-angle annular dark field (HAADF) images were obtained via spherical-aberration corrected field emission TEM (JEM-ARM200FTH, JEOL, Tokyo, Japan) at an accelerating voltage of 200 kV. The dark-field lattice images have a resolution of 0.08 nm. For the AFE and EC behaviors, the current density versus electric field ( $J-E$ ) and the polarization versus electric field ( $P-E$ ) curves of the TiN/ $ZrO_2$ /TiN MIM capacitors with a ZrN interfacial layer were performed using a Keithley 4200-SCS semiconductor analyzer equipped with a 4225-PMU ultrafast IV unit (Tektronix, Inc., Beaverton, OR, USA). A triangular bipolar voltage pulse with a rising and falling time of 50  $\mu s$  was applied to the top electrode of the AFE  $ZrO_2$  capacitors with a ZrN interfacial layer. Meanwhile, the bottom electrode was definitely grounded.

## 3. Results and discussion

### 3.1. Material analyses of TiN/ZrN/ $ZrO_2$ /TiN stacked structures

Figs. 1a and S1 show the XPS depth profiles of atomic percentages of Ti, N, Zr, and O atoms for the TiN/ $ZrO_2$ /TiN stacked structures without and with a ZrN interfacial layer at TiN TE and  $ZrO_2$  interface respectively. From these figures, it can be found that the percentage of oxygen atoms in  $ZrO_2$  thin films became higher with an increase in ALD cycles of the ZrN interfacial layer. Furthermore, the interface between the TiN TE and  $ZrO_2$  thin film was identified to be at the position under an Ar sputtering of approximately 2600 s. Therefore, the O 1s XPS spectra of the TiN/ $ZrO_2$ /TiN stacked structures with different ZrN cycles at the TiN TE and  $ZrO_2$  interface to the middle of  $ZrO_2$  thin films under different Ar sputtering times were investigated and are shown in Figs. 1b and S2a-d respectively. All the spectra can be deconvoluted into three primary peaks: the Zr-O bond at 531.3 eV [25], the Ti-O bond at 530 eV [25], and the non-lattice oxygen at 533.8 eV [26,27]. When the ALD cycles of the ZrN layer increased, the intensities of Ti-O bond reduced owing to the suppression of the interaction between TiN TE and  $ZrO_2$  thin films during the annealing process. Moreover, the peak-to-area ratios (PARs) of Ti-O bond at the TiN TE and  $ZrO_2$  interface to the middle of  $ZrO_2$  thin films under different Ar sputtering times were calculated with respect to the total amount of XPS peak areas and are illustrated in Fig. 1c. The quantified results indicate that the oxidation of TiN TE can be effectively inhibited via the ZrN interfacial layer. The phenomenon can be further confirmed by the Ti 2p XPS spectra (Fig. S3), where Ti-O bonds have been identified at 463 eV ( $2p_{1/2}$ ) and 456.8 eV ( $2p_{3/2}$ ) [28], Ti-N bonds have been determined at 460.8 eV ( $2p_{1/2}$ ) and 455.6 eV ( $2p_{3/2}$ ) [29], and the non-lattice titanium have been detected at 458.4 eV ( $2p_{1/2}$ ) and 454.7 eV ( $2p_{3/2}$ ) [30]. Fig. 1d shows the Zr 3d XPS spectra of the samples with different ZrN cycles at the interface between the TiN TE and  $ZrO_2$  thin films. The spectra can be deconvoluted into three peaks, including Zr-O bonds at 185.3 eV ( $3d_{3/2}$ ) and 183.1 eV ( $3d_{5/2}$ ) [31], Zr-N bonds at 182.2 eV ( $3d_{3/2}$ ) and 180.9 eV ( $3d_{5/2}$ ) [32,33], and the non-lattice zirconium at 181.7 eV ( $3d_{3/2}$ ) and 179.1 eV ( $3d_{5/2}$ ) [34]. Obviously, the peak intensity of Zr-N bonds increased with an increase in ALD cycles of



**Fig. 1.** (a) XPS depth profile of TiN/ZrO<sub>2</sub>/TiN stacked films without a ZrN interfacial layer. (b) O 1s and (d) Zr 3d XPS spectra of TiN/ZrO<sub>2</sub>/TiN stacked films with different ALD cycles of ZrN interfacial layers under a 2600-s Ar sputtering time. (c) XPS PAR of Ti-O bond of TiN/ZrO<sub>2</sub>/TiN stacked films with different ALD cycles of ZrN interfacial layers under different Ar sputtering times.

the ZrN interfacial layer for the oxidation inhibition of TiN TE.

Fig. 2a shows the GIXRD diffractograms of the TiN/ZrO<sub>2</sub>/TiN stacked films with different ZrN cycles at the TiN TE and ZrO<sub>2</sub> interface. According to the diffraction powder data [35], the diffraction peaks, located at 37°, 42.68°, and 61.93° can be identified as the crystallinity of TiN films. Additionally, the peaks located at 30.27°, 35.26°, 50.37°, 59.61°, 60.21°, and 74.54° represent the t(011), t(110), t(112), t(013), t(121), and t(220) peaks of zirconia, respectively, corresponding to the AFE behaviors of ZrO<sub>2</sub> thin films [36]. To further analyze the crystalline structure of the stacked films with different ZrN cycles, the GIXRD spectra in the 2θ range of 27–34° were magnified and are presented in Fig. 2b. In these spectra, the ZrN<sub>9</sub> sample exhibits the highest intensity of t(011) peak. The enhanced crystallinity of ZrO<sub>2</sub> thin films with a ZrN interfacial layer between TiN TE and ZrO<sub>2</sub> can be ascribed to the reduction in interfacial dead layer [37]. Nevertheless, for the ZrN<sub>12</sub> sample, a lower intensity of t(011) peak was observed as compared with that of the ZrN<sub>9</sub>, implying the deficient crystallinity of the tetragonal phase in ZrO<sub>2</sub> thin films owing to the additional stress induced by the thick ZrN interfacial layer [38]. Furthermore, the grain sizes of all samples were extracted to confirm the effectiveness of ZrN interfacial

layer on crystallization of ZrO<sub>2</sub> thin films based on Scherrer equation [39]:

$$\tau = \frac{k\lambda}{\beta \times \cos\theta} \quad (1)$$

where  $\tau$  is the mean size of the crystalline domains, i.e. the grain size of ZrO<sub>2</sub> thin films,  $k$  is the dimensionless shape factor close to unity,  $\lambda$  is the X-ray wavelength of 0.154 nm,  $\beta$  is the full width at half maximum (FWHM) of the peak, and  $\theta$  is the Bragg angle of 30.27°. Fig. 2c presents the calculated grain sizes of the TiN/ZrO<sub>2</sub>/TiN stacked films with different ZrN cycles at the TiN TE and ZrO<sub>2</sub> interface. The grain size of the ZrN<sub>9</sub> sample was close to the film thickness of 10 nm, indicating a complete crystallization of the ZrO<sub>2</sub> thin film. Fig. 3a,b and Fig. S4a show the cross-sectional HRTEM images of the TiN/ZrO<sub>2</sub>/TiN stacked films without and with a ZrN interfacial layer of 9 and 12 cycles, respectively, to visually confirm the crystallinity. Moreover, the crystallinity of the stacked films was further quantified via HAADF imaging, as shown in Fig. 3c,d and Fig. S4b. The bottom figures show the high-magnification images of the lattice with the corresponding fast Fourier

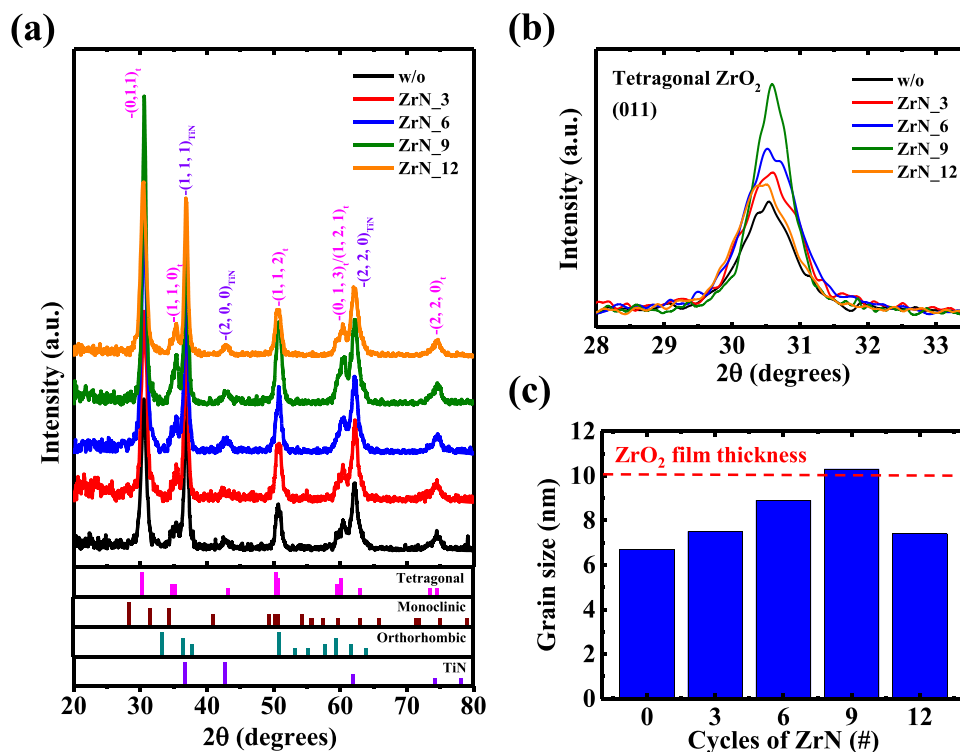


Fig. 2. (a) GIXRD patterns and (b) magnified spectra in the  $2\theta$  range of  $27\text{--}34^\circ$  (tetragonal 011) of TiN/ZrO<sub>2</sub>/TiN stacked films with different ALD cycles of ZrN interfacial layers. The reference numbers of powder diffraction files are TiN (cubic): 01-087-0632, ZrO<sub>2</sub> (tetragonal): 00-050-1089, ZrO<sub>2</sub> (monoclinic): 01-074-1200, and ZrO<sub>2</sub> (orthorhombic): 00-034-1084. (c) Calculated grain sizes of ZrO<sub>2</sub> film films with different ALD cycles of ZrN interfacial layers.

transformed patterns and the simulated atomic columns confirming the tetragonal phase for the matrix. By examining these images, we can observe the highly ordered arrangement of Zr atoms in the ZrN<sub>9</sub> sample, indicating that the optimal cyclic deposition of ZrN interfacial layer can improve the crystallinity of ZrO<sub>2</sub> thin films, corresponding to the calculated grain sizes in Fig. 2c. The enhanced tetragonal phase at the TiN TE and ZrO<sub>2</sub> interface can be attributed to the thinner interfacial dead layer, effectively stabilizing the crystallinity of tetragonal phase at the interface to achieve a large grain. Additionally, the film thickness extracted from the HRTEM images can be used to calculate the electric field of ZrO<sub>2</sub> thin films for further characterizations of AFE and EC behaviors.

### 3.2. Antiferroelectric behaviors of TiN/ZrN/ZrO<sub>2</sub>/TiN capacitors

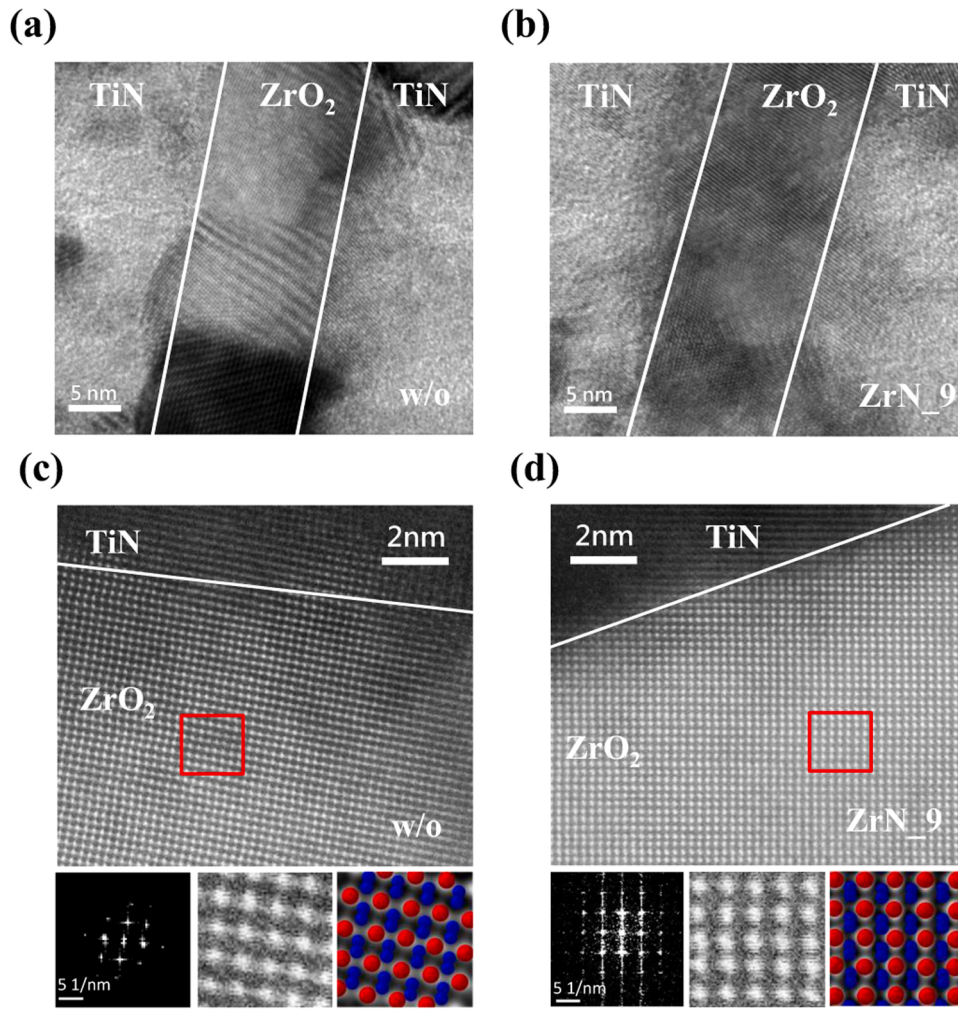
To investigate the effects of the ZrN interfacial layer on the AFE characteristics of TiN/ZrO<sub>2</sub>/TiN MIM capacitors, the  $J\text{--}E$  and  $P\text{--}E$  hysteresis curves were measured using the triangular bipolar pulses with rising and falling times of  $50\ \mu\text{s}$  and are shown in Fig. 4a and b, respectively. According to the film thicknesses obtained from the HRTEM images, the voltages of 5, 5.15, 5.3, 5.45, and 5.6 V, respectively, were applied to w/o, ZrN<sub>3</sub>, ZrN<sub>6</sub>, ZrN<sub>9</sub>, and ZrN<sub>12</sub> samples to get the same electric field of 5 MV/cm. In this figure, the ZrN<sub>9</sub> sample exhibited the largest polarization of  $41.3\ \mu\text{C}/\text{cm}^2$ , corresponding to a more complete crystallization in ZrO<sub>2</sub> thin films, as shown in Fig. 3. Additionally, the critical electric field of the ZrN<sub>9</sub> sample was observed to be significantly reduced. The enhanced AFE behaviors could be attributed to the huge amount of polarized domains in ZrO<sub>2</sub> thin films with a ZrN interfacial layer of 9 cycles [40]. However, the polarization of ZrO<sub>2</sub> thin films with a ZrN interfacial layer of 12 cycles declined obviously. To identify if the polarization can be further enhanced, a larger electric field of 6 MV/cm was applied on the ZrN<sub>12</sub> sample, as shown in Fig. S5. It can be observed that even at a higher electric field, the ZrN<sub>12</sub> sample exhibited a smaller  $P_s$  than that of the ZrN<sub>9</sub>, indicating a poor

crystallinity of ZrO<sub>2</sub> thin films as revealed in Figs. 2 and 3. Fig. 4c shows the energy storage properties of the TiN/ZrO<sub>2</sub>/TiN MIM capacitors with different cycles of ZrN layers at the TiN TE and ZrO<sub>2</sub> interface, where the energy storage density (ESD) and energy loss can be obtained via the integral of the blue and orange areas of the  $P\text{--}E$  hysteresis curves, as illustrated in the inset of Fig. 4b. Furthermore, the efficiency can be calculated based on the following equation:

$$\text{Efficiency} = \frac{\text{ESD}}{\text{ESD} + \text{loss}} \quad (2)$$

The ESD and efficiency of the ZrN<sub>9</sub> sample were significantly enhanced to  $80.2\ \text{J}/\text{cm}^3$  and 64.2 % respectively, suitable for applications in supercapacitors [41], as confirmed by the high dielectric permittivity in the permittivity versus electric field ( $\epsilon\text{--}E$ ) characteristics (Fig. S6). Furthermore, the ESD and efficiency of the samples as a function of the operation frequency was investigated and is shown in Fig. S7. Too high and low frequency operations would degrade the energy storage performance of AFE ZrO<sub>2</sub> thin films because the polarization is screened by the RC delay and leakage current, respectively [42, 43]. Fig. 4d shows the cycling endurance characteristics of the devices stressed at 5 MV/cm with the pulse width of  $50\ \mu\text{s}$  for  $10^6$  cycles to further validate that the ZrN interfacial layer can improve not only the polarization but also reliability behaviors of the TiN/ZrO<sub>2</sub>/TiN stacked films. The  $P\text{--}E$  hysteresis curves of all samples after different numbers of cycling pulses are shown in Fig. S8. When the ALD cycles of ZrN deposition increased, the change in ESD and efficiency of TiN/ZrO<sub>2</sub>/TiN MIM capacitors after  $10^6$  cycling pulses was reduced significantly. It is reported that the degradation of high-k FE and AFE materials after cycling endurance tests was highly correlated with the defects in dielectric thin films [44]. For the devices with a ZrN interfacial layer, the additional nitrogen in ZrN layers can effectively reduce the generation of oxygen vacancies at the TiN TE and ZrO<sub>2</sub> interface during the cycling operations [45,46], enhancing the cycling endurance stability of the TiN/ZrO<sub>2</sub>/TiN MIM capacitors.





**Fig. 3.** HRTEM images of TiN/ZrO<sub>2</sub>/TiN stacked films (a) without and (b) with a ZrN interfacial layer of 9 cycles. HAADF micrographs of TiN/ZrO<sub>2</sub>/TiN stacked films (c) without and (d) with a ZrN interfacial layer of 9 cycles. The bottom figures in (c) and (d) show the high-magnification images of the lattice with the corresponding fast Fourier transformed patterns and the atomic simulation for the tetragonal phase with Zr columns in red and O columns in blue.

In addition to the materials analysis, the pulse switching measurement has been applied to analyze the interface between the FE/AFE films and the metal electrode quantitatively, as proposed by Park et al. [47]. The polarization switching kinetics could be strongly affected by the properties of the interfacial capacitance ( $C_i$ ) originated from the non-polarized layer at the interface [18]. Fig. 5a and b shows the switching current versus time ( $I_{sw} - t$ ) characteristics of the TiN/ZrO<sub>2</sub>/TiN MIM capacitors without and with a ZrN interfacial layer of 9 cycles, respectively. The waveform of voltage applied to the samples during measurement was illustrated in Fig. S9. The  $C_i$  can be extracted from the region 2 of the waveform, representing the switching current of the process from non-polar to polar phases. According to the polarization reversal theory, the switching current can be described as follows [47,48]:

$$I_{sw}(t) = I_{sw}^0 e^{\frac{t-t_0}{R_L C_i}} \quad (t_0 \leq t \leq t_{sw}) \quad (3)$$

where  $t_0$  and  $t_{sw}$  are the times when the film starts switching and when the switching is finished, respectively, and  $R_L$  is the total resistance of the measurement system. Additionally,  $I_{sw}^0$  is the current in the beginning of the polarization switching, which was indicated by the black arrow in these figures and can be expressed as follows:

$$I_{sw}^0 = \frac{(E_a - E_c)t_f}{R_L} \quad (4)$$

where  $E_a$ ,  $E_c$ , and  $t_f$  are the applied electric field, the coercive field, and the thickness of ZrO<sub>2</sub> thin films, respectively.  $R_L$  and  $E_c$  can be extracted from the slope and the intercept of x-axis in  $I_{sw}^0$  versus  $E_a$  characteristics, respectively, as shown in the inset of Fig. 5a and b. Thus, the  $C_i$  can be obtained by substituting the calculated value of  $R_L$  into Eq. (3). Thanks to the modification of the interface between the ZrO<sub>2</sub> thin film and TiN electrode, the  $C_i$  of the ZrN\_9 sample was 43.3  $\mu\text{F}/\text{cm}^2$ , approximately 5.5 times higher than that of the sample without a ZrN interfacial layer, implying that the formation of the interfacial dead layer can be effectively suppressed by the ZrN interfacial layer. To show the improvement intuitively, the schematics of the interface between the ZrO<sub>2</sub> thin films and TiN electrode of the samples without and with a ZrN interfacial layer of 9 cycles are illustrated in Fig. 5c and d, respectively. For the sample without a ZrN interfacial layer, the oxidation of the TiN electrode would occur during the annealing process, causing a thicker interfacial dead layer and fewer domains for the polarization. On the other hand, the ZrN interfacial layer can suppress the reaction between the ZrO<sub>2</sub> thin films and TiN electrode effectively, leading to the improved interface quality and enhanced crystallinity of the ZrO<sub>2</sub> thin films.

### 3.3. Electrocaloric effects of TiN/ZrN/ZrO<sub>2</sub>/TiN capacitors

To investigate the EC effect, the  $P - E$  hysteresis curves of the

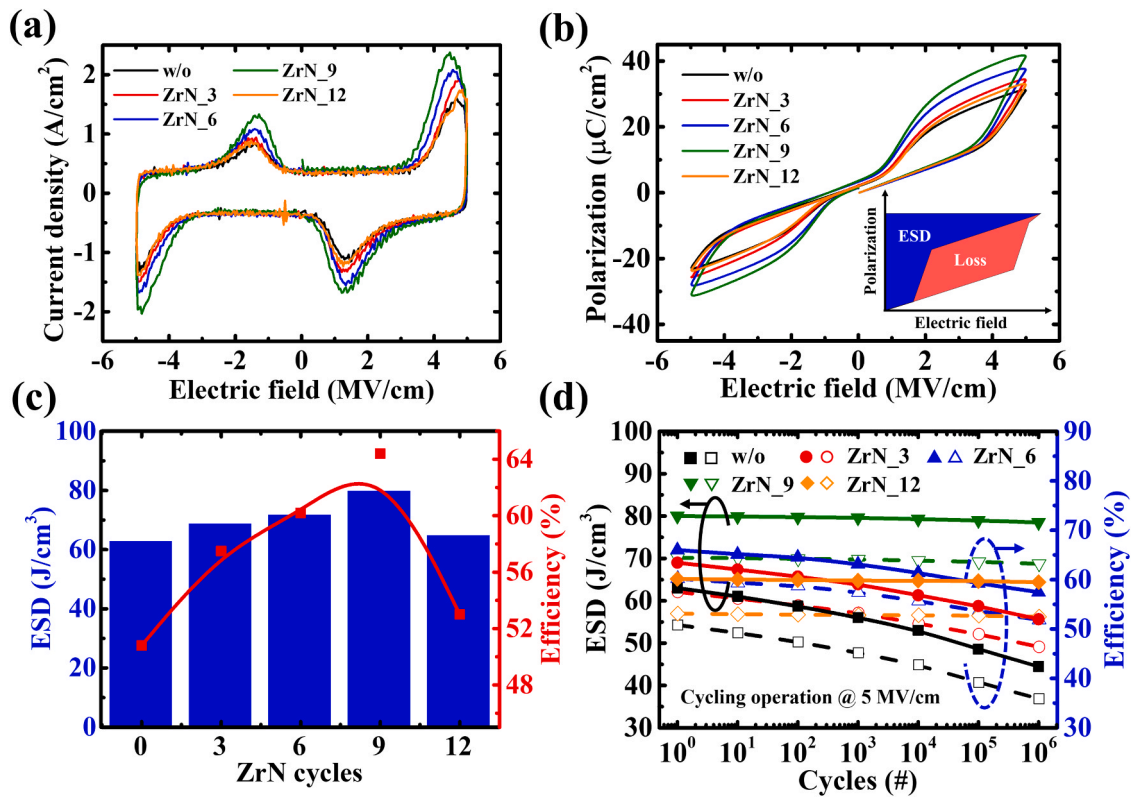


Fig. 4. (a)  $J-E$  and (b)  $P-E$  curves of TiN/ZrO<sub>2</sub>/TiN MIM capacitors with different ALD cycles of ZrN interfacial layers. The inset shows the ESD and loss calculated from the integral of the blue and pink areas of  $P-E$  curves. (c) The values of ESD and efficiency of all samples obtained from  $P-E$  curves. (d) Cycling endurance characteristics of ESD and efficiency for TiN/ZrO<sub>2</sub>/TiN MIM capacitors with different ALD cycles of ZrN interfacial layers.

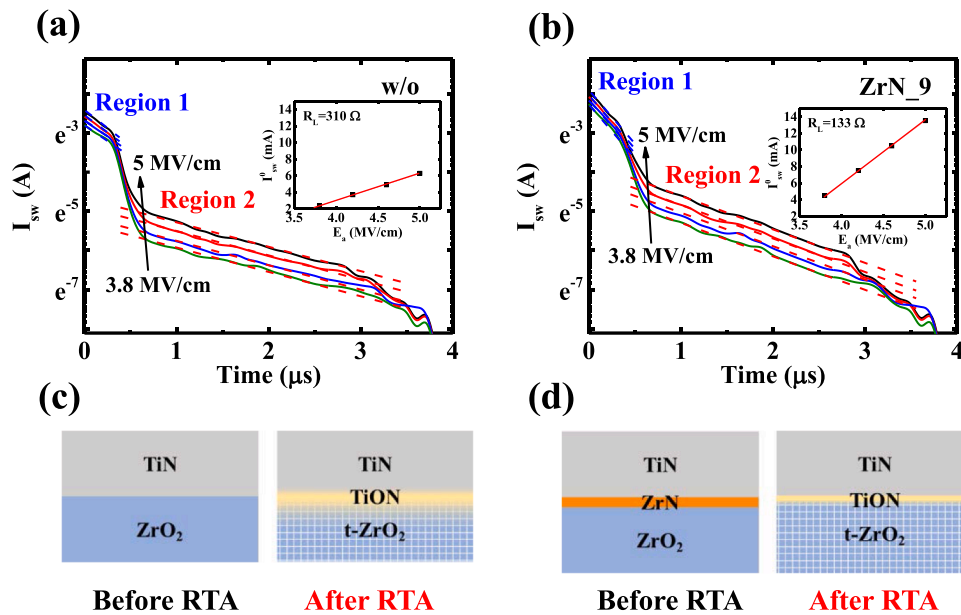


Fig. 5.  $I_{sw}-t$  curves measured at different electric fields of TiN/ZrO<sub>2</sub>/TiN MIM capacitors (a) without and (b) with a ZrN interfacial layer of 9 cycles. The inset shows  $I_{sw}^0-E_a$  curves to extract  $R_L$  and  $E_c$ . Schematics of the interface between the ZrO<sub>2</sub> thin films and TiN electrode of the samples (c) without and (d) with a ZrN interfacial layer of 9 cycles.

samples measured at elevated temperatures were performed. However, the high temperature measurement at high electric field would induce a severe leakage current, as shown in Fig. S10, hindering the EC effect [49]. Thus, the maximum electric field for the polarization measurement was reduced from 5 to 4.5 MV/cm. Fig. 6a and Fig. S11 show the  $P-E$

hysteresis curves of the TiN/ZrO<sub>2</sub>/TiN MIM capacitors with different thicknesses of ZrN interfacial layer measured from 298 K to 388 K to investigate the EC effect. Meanwhile, the temperature dependence on the energy storage performance was examined, as shown in Fig. S12. It can be found that there is no significant trending in energy storage

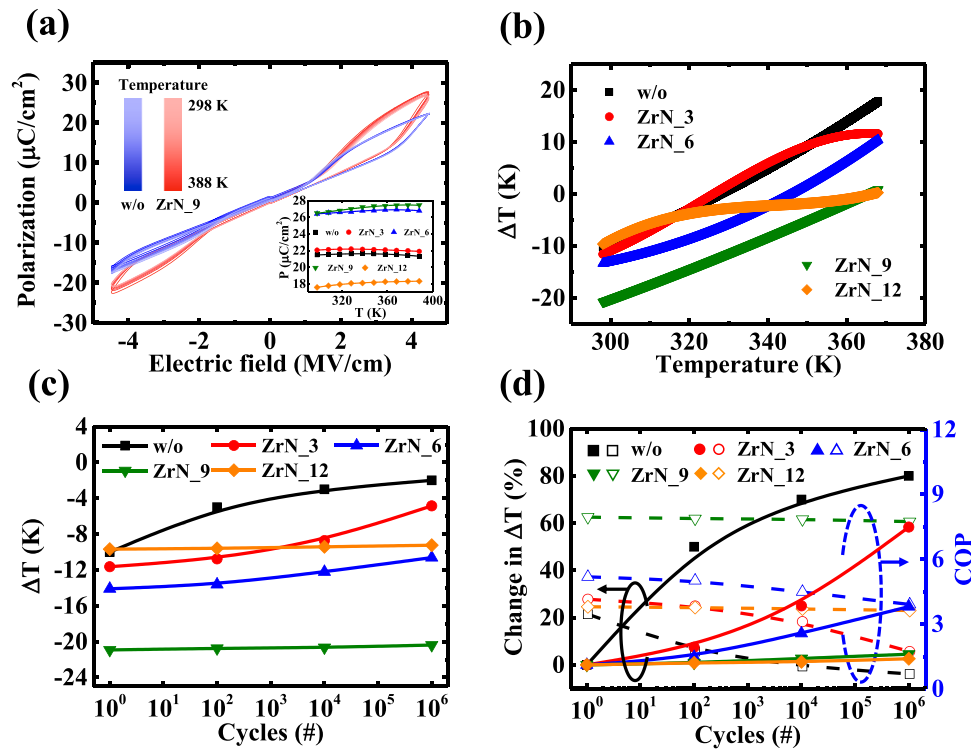


Fig. 6. (a)  $P$ - $E$  curves of TiN/ZrO<sub>2</sub>/TiN MIM capacitors without and with a ZrN interfacial layer of 9 cycles at temperatures ranging from 298 K to 388 K. The insets show the  $P$ - $T$  curves of all samples as the function of temperature estimated from 298 K to 388 K. (b)  $\Delta T$  of all samples as the function of temperature estimated from 298 K to 388 K. (c)  $\Delta T$  and (d) change in  $\Delta T$  and COP of all samples under different numbers of cycling endurance tests.

performance as a function of temperature, suitable for the extraction of adiabatic temperature change of the AFE ZrO<sub>2</sub> thin films. According to the Maxwell's equations, the EC effect can be obtained indirectly via the following formula:

$$\Delta T = -\frac{T}{\rho C_p} \int_{E_1}^{E_2} \left( \frac{\partial P}{\partial T} \right)_E dE \quad (5)$$

where  $\rho$  and  $C_p$  are the density and specific heat capacity of the ZrO<sub>2</sub> thin films, respectively [50,51]. To obtain the numerical differentiation of  $\partial P/\partial T$ , the polarization versus temperature ( $P$  -  $T$ ) characteristics were plotted and fitted using a fourth-order polynomial, as shown in the inset of Fig. 6a. Then, the  $\Delta T$  of all samples at different temperatures was calculated via Eq. (5) and is shown in Fig. 6b. It can be observed that the ZrN<sub>9</sub> sample exhibited a negative EC effect for all the measurement temperatures and an enhanced  $\Delta T$  of  $-21$  K was obtained under 4.5 MV/cm at 298 K, corresponding to the largest polarized response with the change in temperature at room temperature. However, at high temperature, a positive EC effect was observed in the samples with a ZrN interfacial layer of less than 9 cycles because of the low energy barrier for the change between the non-polar and polar phases [52]. Moreover, the cycling endurance tests of the  $\Delta T$  under 4.5 MV/cm at 298 K for 10<sup>6</sup> cycles were performed and are presented in Fig. 6c. The EC effects of all samples measured from 298 K to 388 K after different cycles of endurance test were examined and are showed in Figs. S13–S16. When the thickness of the ZrN interfacial layer increased, the stability of EC effect was significantly improved owing to the suppressed interfacial dead layer as revealed in Fig. 5. To quantify the improvement of the cooling efficiency under the cycling operation, the change in  $\Delta T$  and the coefficient of performance (COP) were calculated and are presented in Fig. 6d. The COP can be calculated via  $|Q|/W_{\text{total}}$ , where  $Q$  is  $T \times \Delta S$  and  $W_{\text{total}}$  is the integral of the electric field with respect to the polarization in  $P$  -  $E$  hysteresis curves of Fig. 6a, Fig. S11, and Figs. S13–S15. Additionally, the  $\Delta S$  is equal to  $\Delta T \times C_p/T$ . For the ZrN<sub>9</sub> sample, the COP was found to be 7.9 and 7.4 before and after the cycling endurance

tests of 10<sup>6</sup> cycles, respectively, which were the largest as compared with others. Fig. 7 shows the benchmarks of the cycling endurance characteristics after the tests of 10<sup>6</sup> cycles in energy storage and EC effect of this study and other literatures [15,53–66]. With a high ESD value of more than 80 J/cm<sup>3</sup>, the changes in ESD and efficiency after the cycling endurance tests of 10<sup>6</sup> cycles of the AFE ZrO<sub>2</sub> capacitors with a ZrN interfacial layer of 9 cycles are comparable to those of other works, as presented in Fig. 7a,b. Furthermore, the highest absolute value of  $\Delta T$  with a low change in  $\Delta T$  after the cycling endurance tests of 10<sup>6</sup> cycles was obtained in this study, as shown in Fig. 7c, indicating that the TiN/ZrO<sub>2</sub>/TiN MIM capacitors with a ZrN interfacial layer are suitable for the device cooling in ICs.

#### 4. Conclusion

In summary, the TiN/ZrO<sub>2</sub>/TiN MIM capacitors with ZrN as the interfacial layer were demonstrated to suppress the redox reaction between the TiN TE and ZrO<sub>2</sub> thin films and enhance the AFE and EC behaviors. With an increase in ALD cycles of the ZrN interfacial layer, the ESD and efficiency of the ZrO<sub>2</sub> capacitors were significantly increased to 80.2 J/cm<sup>3</sup> and 64.2 %, respectively, indicating the more polarized domains in ZrO<sub>2</sub> thin films, as confirmed in the HAADF images. Through the extraction of the interfacial capacitance, the interfacial dead layer of the samples with a ZrN interfacial layer was effectively reduced, corresponding to the less non-polar ZrO<sub>2</sub> at the interface. Furthermore, the cycling endurance stability of energy storage characteristics was improved because of the suppression of oxygen vacancy generation during the cycling operation. The TiN/ZrO<sub>2</sub>/TiN MIM capacitors with a ZrN interfacial layer of 9 cycles exhibited a competitive  $\Delta T$  of  $-21$  K at 298 K and a negligible change in  $\Delta T$  of 4.5 % after a cycling endurance test of 10<sup>6</sup> cycles. Based on the improvement in energy storage and EC behaviors as well as the robust reliability, the ZrN interfacial layer was suggested to be highly promising in AFE ZrO<sub>2</sub>-based devices for energy storage in supercapacitors and solid-state



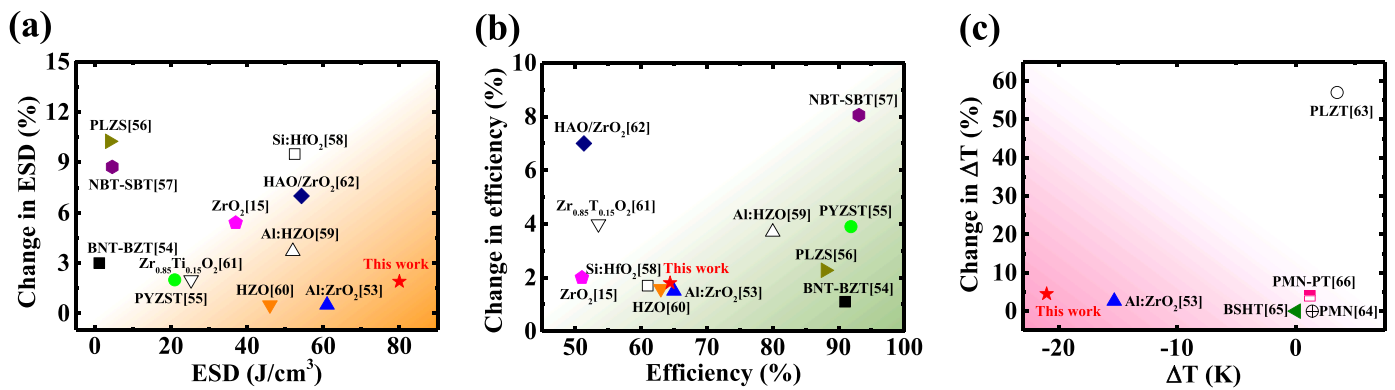


Fig. 7. Benchmarks of cycling endurance characteristics after the tests of  $10^6$  cycles in (a) ESD, (b) efficiency, and (c) EC effect of the TiN/ZrO<sub>2</sub>/TiN MIM capacitor with a ZrN interfacial layer of 9 cycles and other studies.

cooling in ICs.

### CRediT authorship contribution statement

**Yu-Hua Liu:** Investigation, Validation, Data curation, Writing – original draft. **Han-Hsiang Tai:** Investigation, Validation. **Chi-An Ho:** Investigation. **Ting-Han Lin:** Investigation. **Ming-Chung Wu:** Methodology, Validation. **Jer-Chyi Wang:** Conceptualization, Methodology, Writing – review & editing, Supervision.

### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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### Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at [doi:10.1016/j.jeurceramsoc.2023.08.051](https://doi.org/10.1016/j.jeurceramsoc.2023.08.051).

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